



*Simple Derivation of S:* Let  $G(j\omega) = D + C(j\omega I - A)^{-1}B$  be symmetric, noting

$$(j\omega I - A)^{-1} = -A(\omega^2 I + A^2)^{-1} - j\omega(\omega^2 I + A^2)^{-1},$$

The real and imaginary parts of  $G(j\omega)$  are, respectively

$$\operatorname{Re}(G(j\omega)) = D - CA(\omega^2 I + A^2)^{-1}B \quad (8a)$$

$$\operatorname{Im}(G(j\omega)) = -\omega C(\omega^2 I + A^2)^{-1}B. \quad (8b)$$

Similar to the flow of ideas from (2)–(4), a simple exposition for the origin of the singularity matrix  $S$  in [1] can be obtained. To ensure  $\operatorname{Re}(G(j\omega)) > 0, \forall \omega \in \mathbb{R} \cup \infty$ , it is required that  $D = D^T > 0$  and no zero of  $\operatorname{Re}(G(j\omega))$  appear along the positive real  $\omega^2$ -axis. But the zeroes of  $\operatorname{Re}(G(j\omega))$  are given as the poles of its inverse. Recognizing (8a) as a transfer function along  $s = \omega^2$ , its inverse is

$$\left[ \begin{array}{c|c} -A^2 & B \\ \hline -CA & D \end{array} \right]^{-1} = \left[ \begin{array}{c|c} -A^2 + BD^{-1}CA & BD^{-1} \\ \hline D^{-1}CA & D^{-1} \end{array} \right]. \quad (9)$$

Subsequently, the zeros of  $\operatorname{Re}(G(j\omega))$  (i.e., where  $\operatorname{Re}(G(j\omega))$  becomes singular and no longer positive definite), are  $\operatorname{eig}((BD^{-1}C - A)A) = \operatorname{eig}(A(BD^{-1}C - A)) = \operatorname{eig}(S)$ , where  $\operatorname{eig}(\circ)$  denotes the eigenvalues.

**$\operatorname{Re}(G(j\omega)) > 0$  Being a Stronger Passivity Condition:** The discussers present the less obvious result that  $\operatorname{Re}(G(j\omega)) > 0$  (i.e., symmetric positive definite) actually implies  $G(j\omega)$  is symmetric too. To see this, it is noted that a causal and stable physical system must satisfy the Kramers–Krönig relation (see, e.g., [3])

$$\operatorname{Re}(G(j\omega))_{(i,j)} = \frac{1}{\pi} \operatorname{pv} \int_{-\infty}^{+\infty} \frac{\operatorname{Im}(G(j\omega'))_{(i,j)}}{\omega - \omega'} d\omega' \quad (10a)$$

$$\operatorname{Im}(G(j\omega))_{(i,j)} = -\frac{1}{\pi} \operatorname{pv} \int_{-\infty}^{+\infty} \frac{\operatorname{Re}(G(j\omega'))_{(i,j)}}{\omega - \omega'} d\omega' \quad (10b)$$

where “pv” denotes the principal value while the subscript  $(i, j)$  indexes the  $(i, j)$ th element in  $\operatorname{Re}(G(j\omega))$  and  $\operatorname{Im}(G(j\omega))$ . [Equation (10), however, is generally not fulfilled with an unstable transfer function.] That is, the real and imaginary parts of a causal and stable transfer function are not independent and, in particular, a symmetric real part constrains the imaginary part to be symmetric too. This results in  $\operatorname{Re}(G(j\omega)) > 0 \Rightarrow G(j\omega) + G^*(j\omega) = G(j\omega) + G(-j\omega) = 2\operatorname{Re}(G(j\omega)) > 0$ . In other words,  $\operatorname{Re}(G(j\omega)) > 0$  actually serves as a stronger condition for passivity.

The authors’ comments on the aforementioned items would be greatly appreciated.

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**Closure on “A Half-Size Singularity Test Matrix for Fast and Reliable Passivity Assessment of Rational Models”**

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We would like to thank the discussers for their clarification and appreciate their elegant derivation. We already discovered that we were incorrect in claiming the half-size test matrix to be applicable to unsymmetrical models [1]. The error was clarified in a subsequent paper [2]. We would like to mention that we have also derived a half-size test matrix for scattering parameter-based models [3], again applicable only to symmetrical models.

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